

STRENGTH STATISTICAL CHARACTERISTICS OF THE ISOTROPIC MATERIALS WITH DISC-SHAPED DEFECTS

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Abstract. The expressions of the probabilities density distribution of failure loading for an elastic isotropic body under the action of a homogeneous axisymmetric loading are written. The disc-shaped cracks that do not interact with each other are uniformly distributed. The correlation for finding the most probable, mean value, dispersion and variation coefficient of failure loading are obtained. The dependence of the mentioned statistical strength characteristics on the type of applied loading, number of defects (body sizes) and structure material heterogeneity are investigated.

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1. Introduction

The development of the strength theory and criteria of structural materials failure is an important and actual task. Wide application of high-strength materials, which are characterized by brittle fracture, necessitates the creation of reliable methods for predicting the properties of structural elements in the process of exploitation. For effective simulation of brittle materials fracture, it is important to take into account the stochastic nature of their structure. The complex application of known deterministic fracture mechanics solutions and probabilistic statistical methods enables us to construct qualitative models of structural elements strength and reliability for different types of loading. Such a technique has been developed in the articles [1-6].

The purpose of this study is to develop a method for calculating the statistical strength characteristics of brittle model materials under the conditions of homogeneous axisymmetric loading on the basis of defects deterministic theory and probability theory methods.

2. Formulation of the problem

Studying the brittle materials fracture process, the consideration of the defects such as cracks spatial location in a three-dimensional body is important. The simplest model of a three-dimensional body with stochastically distributed defects is a flat, circular-planar (disc-shaped) crack that is characterized by a radius R and two independent orientation angles in space. The model is considerably simplified in the case of an axisymmetric loading. Then, a disc-shaped crack in an axisymmetric stress field is characterized by two statistically independent parameters: the radius R and the orientation angle α between the normal \vec{n} to the crack plane and the symmetry axis Oz .

We shall consider an elastic isotropic body (or its element) that has volume V , and is under the conditions of a homogeneous axisymmetric loading P and $Q = \eta P$. It evenly distributes N disc-shaped cracks that do not interact with each other (the number N will be assumed to be proportional to volume V : $N = N_0 V$, where N_0 is the number of defects per unit volume). The crack resistance of the material K_{IC} will be considered the same throughout the body.

In the case of material isotropy, all the cracks orientations are equally probable ($0 \leq \alpha \leq \pi/2$). Therefore, the end of the normal \vec{n} with the same probability covers the area of the hemisphere $2\pi|\vec{n}|^2$. The ends of the normal with orientation angles α that do not exceed a given value α^* , cover the area of the spherical segment $2\pi|\vec{n}|^2(1 - \cos\alpha)$. The probability that the crack orientation angle α will not exceed a given value α^* is defined as the ratio of the segment areas and the hemisphere: $P = F(\alpha < \alpha^*) = 1 - \cos\alpha$ [7]. Accordingly, the probability distribution density of a random variable α is written as follows: $f(\alpha) = \sin\alpha$.

We choose the probability distribution density of a random parameter R in the form of a generalized β -distribution [8]: $f(R) = \frac{r+1}{d} \left(1 - \frac{R}{d}\right)^r$ ($0 \leq R \leq d$, d is a finite structural characteristic). Then the distribution function is defined as follows: $F(R) = 1 - \left(1 - \frac{R}{d}\right)^{r+1}$. Here r is the parameter of the material fracture (with its increasing the more probable are a small cracks).

3. Distribution function of failure loading

In article [9] the expressions of the failure loading distribution function for a body element with one disc-shaped crack for the following loads are obtained:

1) all-round tension for $P \geq Q \geq 0$ ($0 \leq \eta \leq 1$)

$$F_1(p, \eta) = \begin{cases} \Psi_1(p, \eta) = (1 - p^{-2})^{r+1} - \frac{r+1}{\sqrt{1-\eta}} \int_{p^{-2}}^1 (1-x)^r \sqrt{(p\sqrt{x})^{-1} - \eta} dx, & 1 \leq p \leq 1/\eta; \\ \Psi_2(p, \eta) = (1 - p^{-2})^{r+1} - \frac{r+1}{\sqrt{1-\eta}} \int_{p^{-2}}^{p^{-2}\eta^{-2}} (1-x)^r \sqrt{(p\sqrt{x})^{-1} - \eta} dx, & 1/\eta \leq p < \infty; \end{cases} \quad (1)$$

2) all-round tension for $Q \geq P \geq 0$ ($1 \leq \eta < \infty$)

$$F_1(p, \eta) = \begin{cases} \Psi_3(p, \eta) = \frac{r+1}{\sqrt{1-\eta}} \int_{p^{-2}\eta^{-2}}^1 (1-x)^r \sqrt{(p\sqrt{x})^{-1} - \eta} dx, & 1/\eta \leq p \leq 1; \\ \Psi_4(p, \eta) = (1 - p^{-2})^{r+1} + \frac{r+1}{\sqrt{1-\eta}} \int_{p^{-2}\eta^{-2}}^{p^{-2}} (1-x)^r \sqrt{(p\sqrt{x})^{-1} - \eta} dx, & 1 \leq p < \infty; \end{cases} \quad (2)$$

3) tension in the axial and compression in lateral directions for $P \geq 0, Q \leq 0$ ($-\infty < \eta \leq 0$)

$$F_1(p, \eta) = \Psi_5(p, \eta) = (1 - p^{-2})^{r+1} - \frac{r+1}{\sqrt{1-\eta}} \int_{p^{-2}}^1 (1-x)^r \sqrt{(p\sqrt{x})^{-1} - \eta} dx, \quad (3)$$

$1 \leq p < \infty;$

4) compression in the axial and tension in the lateral directions for $P \leq 0, Q \geq 0$ ($-\infty < \eta \leq 0$)

$$F_1(-p, \eta) = \Psi_6(-p, \eta) = \frac{r+1}{\sqrt{1-\eta}} \int_{p^{-2}\eta^{-2}}^1 (1-x)^r \sqrt{(p\sqrt{x})^{-1} - \eta} dx, \quad -1/\eta \leq -p < \infty. \quad (4)$$

In the expressions (1)-(4) immeasurable loading $p = \frac{2P\sqrt{d}}{\sqrt{\pi}K_{IC}}$ is introduced.

On the basis of the failure loading distribution function (1)-(4), the statistical strength characteristics of the axisymmetric loaded body (or its element) are determined.

Consider partial cases: the same all-round tension ($P = Q$, $\eta = 1$), biaxial tension ($P = 0$, $q = \frac{2Q\sqrt{d}}{\sqrt{\pi}K_{Ic}}$, $\eta \rightarrow \infty$), uniaxial tension ($P > 0$, $Q = 0$, $\eta = 0$), tension in the axial and equal to it compression in lateral directions ($P = -Q > 0$, $\eta = -1$) and compression in the axial and equal to it tension in the lateral directions ($-P = Q > 0$, $\eta = -1$).

Then, the failure loading distribution function for element of the body with one crack will be written as follows:

1) for the same all-round tension

$$F_1(p, 1) = (1 - p^{-2})^{r+1}, \quad 1 \leq p < \infty; \quad (5)$$

2) for biaxial tension

$$F_1(q, \infty) = (r+1) \int_{q^{-2}}^1 (1-x)^r \sqrt{1 - (q\sqrt{x})^{-1}} dx, \quad 1 \leq q < \infty; \quad (6)$$

3) for uniaxial tension

$$F_1(p, 0) = (1 - p^{-2})^{r+1} - (r+1) \int_{p^{-2}}^1 (1-x)^r (p\sqrt{x})^{-1/2} dx, \quad 1 \leq p < \infty; \quad (7)$$

4) for tension in the axial and equal to it compression in lateral directions

$$F_1(p, -1) = (1 - p^{-2})^{r+1} - \frac{r+1}{\sqrt{2}} \int_{p^{-2}}^1 (1-x)^r \sqrt{1 + (p\sqrt{x})^{-1}} dx, \quad 1 \leq p < \infty; \quad (8)$$

5) for compression in the axial and equal to it tension in the lateral directions

$$F_1(-p, -1) = \frac{r+1}{\sqrt{2}} \int_{p^{-2}}^1 (1-x)^r \sqrt{1 + (p\sqrt{x})^{-1}} dx, \quad 1 \leq -p < \infty. \quad (9)$$

4. Distribution density probabilities of failure loading

For a body with a stochastic distribution of N defects, the distribution density probabilities of failure loading are determined [8] as follows:

$$f_N(p, \eta) = N(1 - F_1(p, \eta))^{N-1} \frac{dF_1(p, \eta)}{dp}. \quad (10)$$

Substituting the expressions of the distribution function (1)-(4) into the formula (10), we obtain the expressions of the distribution density probabilities of fracture loading for a body with randomly distributed disc-shaped defect-cracks for the following cases of axisymmetric loading:

1) all-round tension for $P \geq Q \geq 0$ ($0 \leq \eta \leq 1$)

$$f_N(p, \eta) = \begin{cases} N(1 - \Psi_1(p, \eta))^{N-1} \frac{r+1}{2p^{3/2}\sqrt{1-\eta}} \int_{p^{-2}}^1 \frac{(1-x)^r}{\sqrt{\sqrt{x} - \eta px}} dx, & 1 \leq p \leq 1/\eta; \\ (1 - \Psi_2(p, \eta))^{N-1} \frac{r+1}{2p^{3/2}\sqrt{1-\eta}} \int_{p^{-2}}^{p^{-2}\eta^{-2}} \frac{(1-x)^r}{\sqrt{\sqrt{x} - \eta px}} dx, & 1/\eta \leq p < \infty; \end{cases} \quad (11)$$

2) all-round tension for $Q \geq P \geq 0$ ($1 \leq \eta < \infty$)

$$f_N(p, \eta) = \begin{cases} N(1 - \Psi_3(p, \eta))^{N-1} \frac{r+1}{2p^{3/2}\sqrt{1-\eta}} \int_{p^{-2}\eta^{-2}}^1 \frac{(1-x)^r}{\sqrt{\sqrt{x} - \eta px}} dx, & 1/\eta \leq p \leq 1; \\ N(1 - \Psi_4(p, \eta))^{N-1} \frac{r+1}{2p^{3/2}(\eta-1)} \int_{p^{-2}\eta^{-2}}^{p^{-2}} (1-x)^r \sqrt{\frac{1-\eta}{\sqrt{x} - \eta px}} dx, & 1 \leq p < \infty; \end{cases} \quad (12)$$

3) tension in the axial and compression in lateral directions for $P \geq 0, Q \leq 0$ ($-\infty < \eta \leq 0$)

$$f_N(p, \eta) = N(1 - \Psi_5(p, \eta))^{N-1} \frac{r+1}{2p^{3/2}\sqrt{1-\eta}} \int_{p^{-2}}^1 (1-x)^r (\sqrt{x} - \eta px)^{-1/2} dx, \quad 1 \leq p < \infty; \quad (13)$$

4) compression in the axial and tension in the lateral directions for $P \leq 0, Q \geq 0$ ($-\infty < \eta \leq 0$)

$$f_N(p, \eta) = N(1 - \Psi_6(p, \eta))^{N-1} (r+1) \times \left(\frac{2(1-p^{-2}\eta^{-2})^r}{p^3\eta^2} + \frac{1}{2p^{3/2}(\eta-1)} \int_{p^{-2}\eta^{-2}}^1 (1-x)^r \sqrt{\frac{1-\eta}{\sqrt{x} - \eta px}} dx \right), \quad -1/\eta \leq -p < \infty. \quad (14)$$

For partial cases, the distribution density probabilities of failure loading will be written as follows:

1) for the same all-round tension

$$f_N(p,1) = 2N(1 - F_1(p,1))^{N-1} \frac{(r+1)(1-p^{-2})^r}{p^3}, \quad 1 \leq p < \infty; \quad (15)$$

2) for biaxial tension

$$f_N(q,\infty) = \frac{N}{2}(1 - F_1(q,\infty))^{N-1} \frac{r+1}{q^{3/2}} \int_{q^{-2}}^1 \frac{(1-x)^r}{\sqrt{qx - \sqrt{x}}} dx, \quad 1 \leq q < \infty; \quad (16)$$

3) for uniaxial tension

$$f_N(p,0) = \frac{N}{2}(1 - F_1(p,0))^{N-1} \frac{r+1}{p^{3/2}} \int_{p^{-2}}^1 \frac{(1-x)^r}{x^{1/4}} dx, \quad 1 \leq p < \infty; \quad (17)$$

4) for tension in the axial and compression in lateral directions that is equal to it

$$f_N(p,-1) = \frac{N}{2\sqrt{2}}(1 - F_1(p,-1))^{N-1} \frac{r+1}{p^{3/2}} \int_{p^{-2}}^1 \frac{(1-x)^r}{\sqrt{px + \sqrt{x}}} dx, \quad 1 \leq p < \infty; \quad (18)$$

5) for compression in the axial and tension in lateral directions that is equal to it

$$f_N(-p,-1) = N(1 - F_1(-p,-1))^{N-1} (r+1) \left(\frac{2(1-p^{-2})^r}{p^3} - \frac{1}{2\sqrt{2}p^{3/2}} \int_{p^{-2}}^1 \frac{(1-x)^r}{\sqrt{px + \sqrt{x}}} dx \right), \quad 1 \leq -p < \infty. \quad (19)$$

5. Probabilistic characteristics of failure loading

We will investigate some probabilistic characteristics of failure loading. Its most probable value (mode), which corresponds to the loading level, in which the distribution density probabilities reach a maximum, is determined from the equation [8]

$$\frac{d^2 F_N(p,\eta)}{dp^2} = 0, \quad \text{where } F_N(p,\eta) = 1 - (1 - F_1(p,\eta))^N. \quad (20)$$

Equation (20) can be reduced to a form

$$(1-N)(F_1'(p,\eta))^2 + (1 - F_1(p,\eta))F_1''(p,\eta) = 0. \quad (21)$$

Mean value of fracture loading is written as follows [8]:

$$\langle p \rangle = p_{\min}(\eta) + \int_{p_{\min}(\eta)}^{p_{\max}(\eta)} (1 - F_1(p, \eta))^N dp. \quad (22)$$

The dispersion and the coefficient of failure loading variation are found by the formulas [8]

$$D(p) = p_{\min}^2(\eta) + 2 \int_{p_{\min}(\eta)}^{p_{\max}(\eta)} (1 - F_1(p, \eta))^N p dp - \langle p \rangle^2, \quad (23)$$

$$W(p) = \sqrt{D(p)} / \langle p \rangle. \quad (24)$$

Substituting in formulas (20)-(24) the expressions for the failure loading distribution function $F_1(p, \eta)$ (1)-(9), we obtain the relations for determination of the probability strength characteristics considered.

According to the expressions (15)-(19) in Figures 1 and 2 the graphs of the distribution density probabilities of failure loading $f_N(p, \eta)$ for a body with stochastic distribution of N cracks for different types of a stressed state were constructed. Distributions of a failure loading random variable will be unimodal. The threshold value of strength not equal to zero and depends on the type of loading.

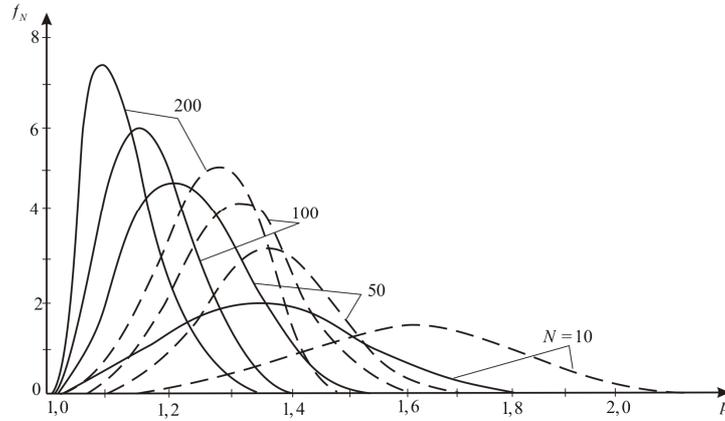


Fig. 1. Distribution density probabilities of failure loading for the same all-round tension for materials with different number of defects and heterogeneity (solid for $r = 3$, dashed for $r = 5$)

According to the expressions (22)-(24) in Figures 3-5, the graphs of failure loading probability characteristics for various types of stressed state for materials with different number of defects and different heterogeneity were constructed.

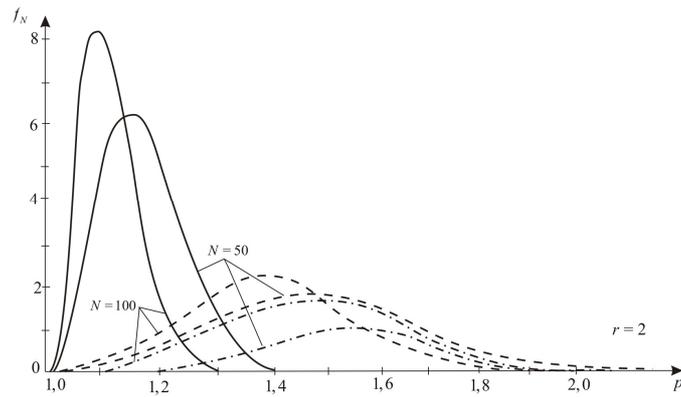


Fig. 2. Distribution density probabilities of failure loading for various types of stressed state (solid for $\eta = 1$, dashed for $\eta = 0$, dotted dashed for $\eta = -1$)

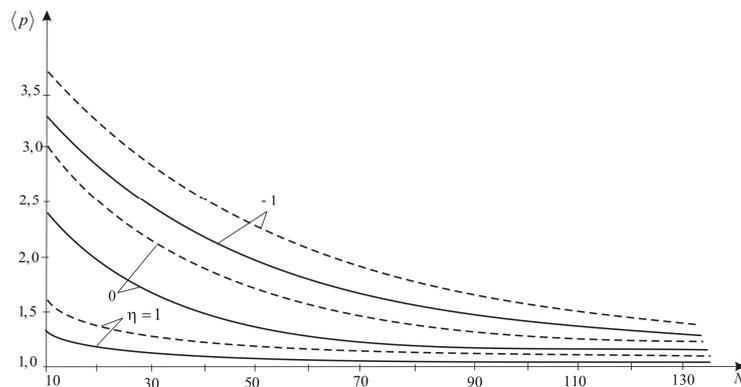


Fig. 3. The failure loading mean value for various types of stressed state (solid for $r = 2$, dashed for $r = 4$)

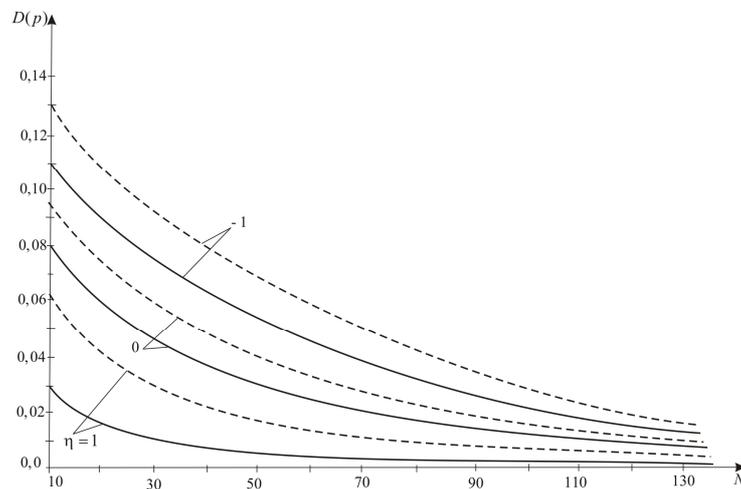


Fig. 4. Dispersion of failure loading for various types of stressed state (solid for $r = 2$, dashed for $r = 4$)

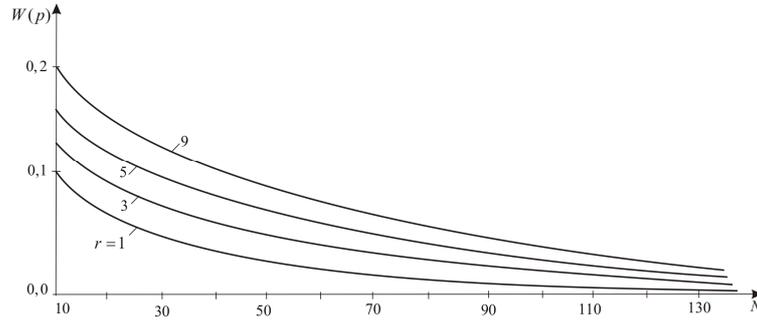


Fig. 5. Coefficient of failure loading variation for materials of different heterogeneity for the same all-round tension

6. Conclusions

In Figures 1 and 2, the influence of cracks number N in the body (body volume at the same defect density), the distribution law of the cracks radiuses, that is, the heterogeneity of the material (parameter r) and the ratio of the applied loading (parameter η) to the most probable value of strength (mode $Mo(p)$) are considered.

As can be seen from Figure 1, for the same all-round tension ($\eta = 1$), with the increasing in the number of cracks in the body, decreases the most probable strength value. This tendency persists when changing material heterogeneity (parameter r). Increasing the value r (increasing the uniformity of the material) leads to an increase $Mo(p)$. With the change of the parameter r , the shape of the distribution density curve changes. The value $Mo(p)$ can also be obtained from equation (21).

In Figure 2, we see a loading type influence on the body strength. The smallest value $Mo(p)$ is observed for the same all-round tension ($\eta = 1$), the greatest for tension in the axial and equal to it compression in lateral directions ($\eta = -1$). Similar conclusions are made in [10], where the indicated values $Mo(p)$ are found in another way. The magnitude $Mo(p)$ increasing and shape of the distribution density curve are affected by the number of defects N .

We note that the maximum ordinates of the distribution curve are directly proportional to N and inversely proportional to r . This feature does not depend on the type of applied loading. Consequently, in the case of N increasing, the maximum value of the distribution density also increases and decreases with r increasing.

In Figure 3, the influence of the loading ratio, the number of cracks and material homogeneity on the failure loading mean value $\langle p \rangle$ is considered. With material homogeneity increasing, the failure loading mean value increases. Its greatest value will be in case of tension in the axial and equal to it compression in lateral directions. Such regularities are observed in [9]. Note that there is a certain range of body sizes, for which strength with an asymptotic approximation to its threshold value almost does not depend on the defects number.

Figure 4 shows the dependence of the failure loading dispersion $D(p)$ on the number of cracks, the loading ratio and the material homogeneity. The dispersion is a decreasing function of the number of cracks in the body. At a certain interval of N change we see a rapid decreasing of $D(p)$ value. The nature of this decreasing does not depend on the type of loading and parameter r . With a change in the parameter r , the strength dispersion varies in size, which is almost independent of η . As in the case of the mean value $\langle p \rangle$, there is a certain range of body sizes, for which the strength dispersion is almost independent of the defects number.

In Figure 5, the influence of material homogeneity and the number of cracks on the coefficient of failure loading variation $W(p)$ for the same all-round tension is investigated. The magnitude $W(p)$ increases with the parameter r increasing and decreases with the cracks number N increasing. There is a certain range of body sizes for which we observe a significant change in the $W(p)$ magnitude and an asymptotic approximation to a certain threshold value. Similar patterns are traced to other types of a stressed state.

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