# METHOD OF FUNDAMENTAL SOLUTIONS FOR HELMHOLTZ EIGENVALUE PROBLEMS IN ELLIPTICAL DOMAINS 

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#### Abstract

In this paper, the method of fundamental solutions for Helmholtz eigenproblems in an elliptical domain is presented. To find the approximate solution of the problem, the Hankel function of the first kind and zero order as the fundamental solution of the Helmholtz equation in unbounded domain on the plane was used. Numerical examples illustrating the accuracy of the present method are given.


## Introduction

The transverse vibration of a membrane which occupy the domain $S$, is governed by the following equation [1]

$$
\begin{equation*}
\nabla^{2} w=\frac{1}{c^{2}} \frac{\partial^{2} w}{\partial t^{2}} \tag{1}
\end{equation*}
$$

where $\nabla^{2}$ is the Laplace operator, $c=\sqrt{T / \rho}$ is the speed of sound, $T$ is the uniform tension per unit length of the edge $\partial S, \rho$ is the density per unit area, $w(x, y, t)$ is the displacement of the plate point $(x, y)$ at time $t$. When the free vibration of the membrane is considered, a harmonic time dependence of the form $e^{i \omega t}$ is assumed, i.e. $w(x, y, t)=W(x, y) e^{i \omega t}$. Substituting the form of the displacement into (1), the Helmholtz equation is obtained

$$
\begin{equation*}
\nabla^{2} W+\Omega^{2} W=0, \quad(x, y) \in S \tag{2}
\end{equation*}
$$

where $\Omega=\omega \sqrt{\rho / T}$. The equation (2) is completed by boundary conditions. We assume here the Dirichlet condition

$$
\begin{equation*}
W(x, y)=0, \quad(x, y) \in \partial S \tag{3}
\end{equation*}
$$

The differential equation (2) and boundary condition (3) form the Helmholtz eigenvalue problem. The problem for elliptic domain $S$ is the subject of this paper.

## 1. The method of fundamental solutions

We solve the eigenvalue problem (2)-(3) by using the method of fundamental solutions (MFS). The MFS is a boundary method which does not involve discretization and integration. The method is applicable to problems for which a fundamental solution of the governing equation is known. The basic idea of the method is the usage of a linear combination of fundamental solutions with sources located at fictitious points outside the domain of the considered problem.

The fundamental solution of a differential equation is not unique. For the Helmholtz equation (2) we take the Hankel function [2, 3]

$$
\begin{equation*}
G(P, Q)=\frac{i}{4} H_{0}^{(1)}(\Omega d(P, Q)) \tag{4}
\end{equation*}
$$

where $H_{0}^{(1)}(\cdot)$ is the Hankel function of the first kind and zero order, $d(P, Q)$ is the distance between points $P$ and $Q, i=\sqrt{-1}$. The functions $G\left(x, y ; Q_{k}\right)$ satisfy the Helmholtz equation in the domain $S$ for each source points $Q_{k}$ located outside $S$. In the MFS, we approximate the solution of the problem by a function of the form

$$
\begin{equation*}
w_{n}(P)=\sum_{k=1}^{n} c_{k} G\left(P, Q_{k}\right) \tag{5}
\end{equation*}
$$

The approximate solution $w_{n}$ satisfies the differential equation (2), and it does not satisfy the boundary condition (3). The condition can be satisfied approximately by a suitable determination of the coefficients $c_{k}, k=1,2, \ldots, n$. For this purpose we use the last square method. First we choose the points $P_{i}, i=1,2, \ldots, n$, located on boundary $\partial S$ of the domain $S$. We determine the coefficients $c_{k}$ so that the function

$$
\begin{equation*}
f\left(\left(c_{1}, c_{2}, \ldots, c_{n}\right)\right)=\sum_{i=1}^{n}\left[w_{n}\left(P_{i} ; c_{1}, c_{2}, \ldots, c_{n}\right)\right]^{2} \tag{6}
\end{equation*}
$$

has a minimum. This leads to a homogeneous linear system of equations

$$
\begin{equation*}
\frac{\partial f}{\partial c_{j}}=2 \sum_{k=1}^{n} c_{k} \sum_{i=1}^{n} G\left(P_{i}, Q_{j}\right) G\left(P_{i}, Q_{k}\right)=0, \quad j=1,2, \ldots, n \tag{7}
\end{equation*}
$$

which can be written in the matrix form as

$$
\begin{equation*}
\mathbf{A} \mathbf{c}=\mathbf{0} \tag{8}
\end{equation*}
$$

where $\mathbf{A}=\left[a_{j k}\right]_{1 \leq j, k \leq n}, a_{j k}=\sum_{i=1}^{n} G\left(P_{i}, Q_{j}\right) G\left(P_{i}, Q_{k}\right), c=\left[\begin{array}{llll}c_{1} & c_{2} & \ldots & c_{n}\end{array}\right]^{T}$

For a non-trivial solution of equation (8), the determinant of the coefficient matrix $\mathbf{A}$ is set equal to zero, yielding the eigenvalue equation

$$
\begin{equation*}
\operatorname{det} \mathbf{A}(\Omega)=0 \tag{9}
\end{equation*}
$$

Equation (9) with the unknown $\Omega$, is solved numerically. The eigenfunctions for the determined eigenvalues $\Omega_{m}, m=1,2, \ldots$, are given by (5) where the coefficients $c_{k}, k=2, \ldots, n$, are derived in dependence of $c_{1}$ from $n-1$ equations of the system (7).

For the considered Helmholtz eigenvalue problem the function $w_{n}$ defined by (5), after using (4) has the form

$$
\begin{equation*}
w_{n}(x, y)=\sum_{k=1}^{n} c_{k} H_{0}^{(1)}\left(\Omega \sqrt{\left(x-\xi_{k}\right)^{2}+\left(y-\eta_{k}\right)^{2}}\right), \quad(x, y) \in S \tag{10}
\end{equation*}
$$

where the source points $Q_{k}\left(\xi_{k}, \eta_{k}\right), k=1,2, \ldots, n$, are over the domain $S$. But the coefficients of the matrix $\mathbf{A}$ are

$$
\begin{equation*}
a_{j k}=\sum_{i=1}^{n} H_{0}^{(1)}\left(\Omega \sqrt{\left(x_{i}-\xi_{j}\right)^{2}+\left(y_{i}-\eta_{j}\right)^{2}}\right) H_{0}^{(1)}\left(\Omega \sqrt{\left(x_{i}-\xi_{k}\right)^{2}+\left(y_{i}-\eta_{k}\right)^{2}}\right) \tag{11}
\end{equation*}
$$

where the collocation points $P\left(x_{i}, y_{i}\right), i=1,2, \ldots, n$, are located on the boundary $\partial S$ of the domain $S$.

## 2. Numerical examples

We assume that the boundary of the considered membrane domain is an ellipse with major and minor semi-diameters equal $a$ and $b$, respectively. The collocation points are uniformly distributed on this ellipse and the source points are uniformly distributed on a circle with the radius $r=\alpha \cdot \max \{a, b\}, \alpha>1$. Geometry configuration of an elliptical domain with the collocation points on the boundary and a circle with the source points is presented in Figure 1.
The Hankel function $H_{0}^{(1)}$ occurring in equations $(10,11)$ assumes complex values and that way the elements of the matrix $\mathbf{A}$ are complex. Our goal is to evaluate real roots of the complex equation (9). For this purpose the new function $F$ is defined as

$$
\begin{equation*}
F(\Omega)=|\operatorname{det} \mathbf{A}(\Omega)| \tag{12}
\end{equation*}
$$

where the symbol $|\cdot|$ denotes a modulus of complex number. It is clear that the roots of equation (9) are determined by minima of the function $F$. In Figure 2 an
exemplary graph of the function in logarithmic scale for an elliptical domain with semi-diameters ratio $a / b=2.0$ and $n=24$ is presented. Locations of the distinct drops indicate the roots of equation (9).


Fig. 1. Geometry configuration of the considered domain with collocation points $P_{i}$ on the ellipse and source points on $Q_{j}$ the circle


Fig. 2. Logarithmic values of modulus of the determinant $\operatorname{det} \mathbf{A}(\Omega)$ as a function of $\Omega$ for the Helmholtz operator in the elliptical domain with semi-diameters ratio $a / b=2.0$

To estimate the MFS accuracy, an example of the Helmholtz eigenvalue problem in a circular domain will be solved. In this case the eigenvalues are the roots of the equation [3]

$$
\begin{equation*}
J_{m}(\Omega)=0, m=1,2, \ldots \tag{13}
\end{equation*}
$$

A sequence of the eigenvalues for each $m$ can be obtained. First members of the sequences are numerically determined by finding the local minima of the function $\log |F(\Omega)|$, where $F(\Omega)$ is given by equation (12). The calculations were performed for various numbers of source points ( $n=16 ; 20 ; 24$ ). The comparison of the results obtained by the MFS with exact values obtained as roots of equation (13) is presented in Table 1. Note that the numerical calculation by MFS with small number of the source points gives the results with a larger error and can lead to omission of eigenvalues. The comparison of the results obtained by MFS and the exact eigenvalues shown that the absolute error for $n=24$ is not larger as $10^{-5}$.

The eigenvalues of the Helmholtz operator are the frequency parameters of free vibration of a membrane. The first seven frequency parameter values of the elliptic membrane with clamped edge are presented in Table 2. The calculations are performed for various values of semi-diameters ratio $a / b$ of the ellipse. The frequency parameters obtained by using MFS were compared with the eigenvalues determined by the finite element method in the paper [1] by Buchanan and Peddieson. For assumed number of sources $n=24$ in MFS a small differences of the results calculated by using both methods are observed.

Table 1. Eigenvalues $\Omega_{n}(n=1, \ldots, 10)$ of the Helmholtz operator in a circular domain obtained by the MFS for $n=16 ; 20 ; 24$ and exact eigenvalues obtained as roots of equation (12)

|  | MFS <br> $n=16$ | MFS <br> $n=20$ | MFS <br> $n=24$ | Exact <br> eigenvalues |
| :---: | :---: | :---: | :---: | :---: |
| $\Omega_{1}$ | 2.40483 | 2.40483 | 2.40483 | 2.40483 |
| $\Omega_{2}$ | 3.83172 | 3.83171 | 3.83171 | 3.83171 |
| $\Omega_{3}$ | 5.13560 | 5.13562 | 5.13562 | 5.13562 |
| $\Omega_{4}$ | 5.52005 | 5.52008 | 5.52008 | 5.52008 |
| $\Omega_{5}$ | 6.37893 | 6.38014 | 6.38016 | 6.38016 |
| $\Omega_{6}$ | 7.01522 | 7.01558 | 7.01559 | 7.01559 |
| $\Omega_{7}$ | 7.58199 | 7.58809 | 7.58834 | 7.58834 |
| $\Omega_{8}$ | 8.40860 | 8.41726 | 8.41724 | 8.41724 |
| $\Omega_{9}$ | 8.65317 | 8.65374 | 8.65373 | 8.65373 |
| $\Omega_{10}$ | - | 8.76646 | 8.77148 | 8.77148 |

Table 2. Frequency parameters $\Omega_{n}(n=1, \ldots, 7)$ of the elliptic membrane for various values of semi-diameters ratio $b / a$ of the ellipse

|  | $\mathrm{a} / \mathrm{b}=1.5$ |  | $\mathrm{a} / \mathrm{b}=2.0$ |  | $\mathrm{a} / \mathrm{b}=3.0$ |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | MFS | FEM $^{*}$ | MFS | FEM $^{*}$ | MFS | FEM $^{*}$ |
| $\Omega_{1}$ | 2.03911 | 2.039 | 1.88858 | 1.889 | 1.76311 | 1.763 |
| $\Omega_{2}$ | 2.91562 | 2.917 | 2.50508 | 2.508 | 2.14359 | 2.148 |
| $\Omega_{3}$ | 3.84097 | 3.850 | 3.42590 | 3.427 | 2.98058 | 3.072 |
| $\Omega_{4}$ | 4.78926 | 4.828 | 3.99047 | 4.000 | 3.42586 | 3.477 |
| $\Omega_{5}$ | 5.10279 | 5.104 | 4.58516 | 4.584 | 3.68419 | 3.697 |
| $\Omega_{6}$ | 5.74644 | 5.769 | 4.98843 | 4.990 | 4.04413 | 4.130 |
| $\Omega_{7}$ | 6.02787 | 6.127 | 5.54225 | 5.561 | 4.84154 | 4.821 |

the results by FEM are given in the paper [4]

## Conclusions

In the paper, the application of the method of fundamental solutions to the Helmholtz eigenvalue problem in the elliptical domain has been presented. As the fundamental solution of the Helmholtz equation in the plane, the Hankel function of the first kind and zero order was used. In order to determine the eigenvalues, the minimum of a real function which is the logarithm of modulus of a complex function was found. The source points occurring in the approximate formula of the solution can be selected on a circle outside of the considered elliptical domain. The comparison of numerical results shows that high accuracy of the calculation is achieved for 24 sources.

## References

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