

## FRACTIONAL HEAT CONDUCTION IN MULTILAYER SPHERICAL BODIES

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**Abstract.** In this paper an analytical solution of the time-fractional heat conduction problem in a spherical coordinate system is presented. The considerations deal the two-dimensional problem in multilayer spherical bodies including a hollow sphere, hemisphere and spherical wedge. The mathematical Robin conditions are assumed. The solution is a sum of time-dependent function satisfied homogenous boundary conditions and of a solution of the steady-state problem. Numerical example shows the temperature distributions in the hemisphere for various order of time-derivative.

**Keywords:** *heat conduction, multilayer bodies, Caputo derivative, spherical coordinate*

### 1. Introduction

The heat conduction in multilayer bodies with assumption of the Fourier law of heat transfer has been considered by Özişik in book [1]. Derivations of the temperature distributions in the bodies in rectangular, cylindrical and spherical coordinate systems were presented. The heat conduction in layered spheres with time-dependent boundary conditions assuming spherical symmetry was the subject of paper [2]. A solution of the heat conduction problem for a two-dimensional multi-layered sphere, hemisphere, spherical cone and spherical wedge presents paper [3].

A generalization of the Fourier law leads to a fractional heat conduction equation. The fractional differential equation governing the heat equation includes the fractional derivatives with respect to time and/or space variables. Properties of the fractional derivatives and methods to the solution of the fractional differential equations are presented in books [4-6]. A method of solution of a time-fractional heat conduction equation in a solid sphere has been discussed by Ning and Jiang in paper [7]. The time-fractional heat conduction in a multi-layered solid sphere assuming spherical symmetry was the subject of paper [8]. Heat conduction modelling using fractional order derivatives is presented by Žecová and J. Terpák in paper [9].

In this paper, an analytical solution of the time-fractional heat conduction for two-dimensional multilayered spherical bodies is presented. The condition for ensuring the perfect contact at interfaces and the mathematical Robin boundary conditions at boundary surfaces are assumed.

## 2. Formulation of the problem

Consider  $n$  spherical concentric layers which are in perfect thermal contact. The  $i$ -th layer ( $i=1, \dots, n$ ) occupied a region described by the spherical coordinates:  $r_{i-1} \leq r \leq r_i$ ,  $0 \leq \varphi \leq \bar{\varphi}$  ( $0 < \bar{\varphi} \leq \pi$ ),  $0 \leq \phi < 2\pi$ , where  $r$  is the radial coordinate,  $\varphi$  is the polar angle and  $\phi$  is the azimuthal angle (Fig. 1).

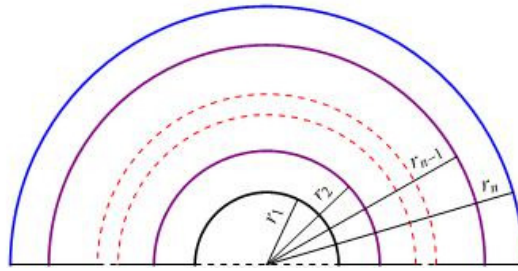


Fig. 1. A schematic diagram of the  $n$ -layered hemisphere

We suppose that the  $i$ -th layer is characterized by constant thermal conductivity  $\lambda_i$  and constant thermal diffusivity  $a_i$ . Moreover, assuming that the temperature doesn't depend on the azimuthal angle, the time fractional heat conduction in the layers is governed by the following differential equation [10]:

$$\frac{1}{r^2} \frac{\partial}{\partial r} \left( r^2 \frac{\partial T_i}{\partial r} \right) + \frac{1}{r^2 \sin \varphi} \frac{\partial}{\partial \varphi} \left( \sin \varphi \frac{\partial T_i}{\partial \varphi} \right) = \frac{1}{a_i} \frac{\partial^\alpha T_i}{\partial t^\alpha} \quad (1)$$

where  $\alpha$  denotes fractional order of the Caputo derivative with respect to time  $t$ ,  $0 < \alpha \leq 1$ . The Caputo derivative is defined by [11]

$${}_0^C D_t^\alpha f(t) = \frac{d^\alpha f(t)}{dt^\alpha} = \frac{1}{\Gamma(m-\alpha)} \int_0^t (t-\tau)^{m-\alpha-1} \frac{d^m f(\tau)}{d\tau^m} d\tau, \quad m-1 < \alpha < m \quad (2)$$

In order to simplify the equation (1) we introduce a new variable  $\mu$  which is related to the polar angle  $\varphi$  by

$$\mu = \cos \varphi \quad (3)$$

Taking into account this relationship in equation (1) we obtain

$$\frac{1}{r^2} \frac{\partial}{\partial r} \left( r^2 \frac{\partial}{\partial r} T_i(r, \mu, t) \right) + \frac{1}{r^2} \frac{\partial}{\partial \mu} \left( (1 - \mu^2) \frac{\partial}{\partial \mu} T_i(r, \mu, t) \right) = \frac{1}{a_i} \frac{\partial^\alpha}{\partial t^\alpha} (T_i(r, \mu, t)), \quad (4)$$

$$i = 1, \dots, n, \quad \mu_\varphi \leq \mu \leq 1$$

where  $\mu_\varphi = \cos \bar{\varphi}$ . The differential equations (4) are complemented by boundary conditions and the conditions providing the perfect thermal contact of the neighbouring layers. The mathematical conditions are [2, 10]

$$\lambda_1 \frac{\partial T_1}{\partial r}(r_0, \mu, t) = -a_+ (T_+(\mu) - T_1(r_0, \mu, t)), \quad (5)$$

$$-\lambda_n \frac{\partial T_n}{\partial r}(r_n, \mu, t) = -a_\infty (T_\infty(\mu) - T_n(r_n, \mu, t)), \quad (6)$$

$$T_i(r, \mu_\varphi, t) = 0 \quad \text{or} \quad \frac{\partial T_i}{\partial \mu}(r, \mu_\varphi, t) = 0 \quad (7)$$

$$T_i(r_i, \mu, t) = T_{i+1}(r_i, \mu, t), \quad i = 1, \dots, n-1 \quad (8)$$

$$\lambda_i \frac{\partial T_i}{\partial r}(r_i, \mu, t) = \lambda_{i+1} \frac{\partial T_{i+1}}{\partial r}(r_i, \mu, t), \quad i = 1, \dots, n-1 \quad (9)$$

where  $a_+$ ,  $a_\infty$  are inner and outer heat transfer coefficients and  $T_+$ ,  $T_\infty$  are inner and outer ambient temperatures. The initial condition is assumed in the form

$$T_i(r, \mu, 0) = F_i(r, \mu) \quad (10)$$

### 3. Solution of the problem

We search a solution to the problem (1), (4)-(10) in the form of a sum

$$T_i(r, \mu, t) = \theta_i(r, \mu, t) + \Phi_i(r, \mu), \quad i = 1, 2, \dots, n \quad (11)$$

where the function  $\theta_i(r, \mu, t)$  satisfies homogeneous fractional heat conduction differential equation and homogeneous boundary conditions and the function  $\Phi_i(r, \mu)$  is a solution of a steady-state problem. Substituting the solution (11) into equations (1), (4)-(10) we obtain the problems for the functions  $\theta_i$  and  $\Phi_i$ . For  $\theta_i$  we have

$$\frac{1}{r^2} \frac{\partial}{\partial r} \left( r^2 \frac{\partial}{\partial r} \theta_i(r, \mu, t) \right) + \frac{1}{r^2} \frac{\partial}{\partial \mu} \left( (1 - \mu^2) \frac{\partial}{\partial \mu} \theta_i(r, \mu, t) \right) = \frac{1}{a_i} \frac{\partial^\alpha}{\partial t^\alpha} (\theta_i(r, \mu, t)), \quad i = 1, \dots, n \quad (12)$$

$$\lambda_1 \frac{\partial \theta_1}{\partial r}(r_0, \mu, t) = a_+ \theta_1(r_0, \mu, t) \quad (13)$$

$$-\lambda_n \frac{\partial \theta_n}{\partial r}(r_n, \mu, t) = a_\infty \theta_n(r_n, \mu, t) \quad (14)$$

$$\theta_i(r, \mu_\varphi, t) = 0 \quad \text{or} \quad \frac{\partial \theta_i}{\partial \mu}(r, \mu_\varphi, t) = 0 \quad (15)$$

$$\theta_i(r_i, \mu, t) = \theta_{i+1}(r_i, \mu, t), \quad i = 1, \dots, n-1 \quad (16)$$

$$\lambda_i \frac{\partial \theta_i}{\partial r}(r_i, \mu, t) = \lambda_{i+1} \frac{\partial \theta_{i+1}}{\partial r}(r_i, \mu, t), \quad i = 1, \dots, n-1 \quad (17)$$

$$\theta_i(r, \mu, 0) = F_i(r, \mu) - \Phi_i(r, \mu) = F_i^*(r, \mu) \quad (18)$$

The functions  $\Phi_i(r, \mu)$  satisfy the following boundary problem

$$\frac{\partial}{\partial r} \left( r^2 \frac{\partial}{\partial r} \Phi_i(r, \mu) \right) + \frac{\partial}{\partial \mu} \left( (1 - \mu^2) \frac{\partial}{\partial \mu} \Phi_i(r, \mu) \right) = 0 \quad (19)$$

$$\lambda_1 \frac{\partial \Phi_1}{\partial r}(r_0, \mu) = -a_+ (T_+(\mu) - \Phi_1(r_0, \mu)), \quad (20)$$

$$-\lambda_n \frac{\partial \Phi_n}{\partial r}(r_n, \mu) = -a_\infty (T_\infty(\mu) - \Phi_n(r_n, \mu)), \quad (21)$$

$$\Phi_i(r, \mu_\varphi) = 0 \quad \text{or} \quad \frac{\partial \Phi_i}{\partial \mu}(r, \mu_\varphi) = 0 \quad (22)$$

$$\Phi_i(r_i, \mu) = \Phi_{i+1}(r_i, \mu), \quad i = 1, \dots, n-1 \quad (23)$$

$$\lambda_i \frac{\partial \Phi_i}{\partial r}(r_i, \mu) = \lambda_{i+1} \frac{\partial \Phi_{i+1}}{\partial r}(r_i, \mu), \quad i = 1, \dots, n-1 \quad (24)$$

### 3.1. Solution to the homogeneous problem

We find the time dependent function  $\theta_i$  as a solution of the problem (12)-(18), by using the separation of variables method. Substituting the product of functions

$$\theta_i(r, \mu, t) = R_i(r) \Psi(\mu) \Lambda(t) \quad (25)$$

into equation (12), we obtain three differential equations

$$\frac{d}{d\mu} \left( (1-\mu^2) \frac{d\Psi(\mu)}{d\mu} \right) = -\nu \Psi(\mu) \quad (26)$$

$$\frac{1}{r^2} \frac{d}{dr} \left( r^2 \frac{d}{dr} R_i(r) \right) + \left[ \omega_i^2 - \frac{\nu}{r^2} \right] R_i(r) = 0 \quad (27)$$

$$\frac{d^\alpha}{dt^\alpha} \Lambda(t) + \gamma^2 \Lambda(t) = 0 \quad (28)$$

where  $\omega_i = \frac{\gamma}{\sqrt{a_i}}$ ,  $\gamma$  and  $\nu$  are separation constants.

#### *Solution of the equation (26)*

Assuming that  $\nu = \beta(\beta+1)$ , the general solution of equation (26) can be written in the form

$$\Psi(\mu) = A_1 P_\beta(\mu) + A_2 Q_\beta(\mu) \quad (29)$$

where  $P_\beta(\mu)$  and  $Q_\beta(\mu)$  are the Legendre functions of order  $\beta$ . Because  $\lim_{\mu \rightarrow 1} Q_\beta(\mu) = \infty$ , we assume  $A_2 = 0$ . Taking into account that

$$\Psi(\mu) = A_1 P_\beta(\mu) \quad (30)$$

and substituting function  $\theta_i(r, \mu, t)$  in the form (25) into equation (15), we obtain an eigenvalue equation

$$P_\beta(\mu_\varphi) = 0 \quad \text{or} \quad \mu_\varphi P_\beta(\mu_\varphi) - P_{\beta+1}(\mu_\varphi) = 0 \quad (31)$$

Solving equation (31a) or (31b), we obtain a sequence of roots  $(\beta_m)_{m=1,2,\dots}$ . The functions  $P_{\beta_m}(\mu)$  corresponding to the values of  $\beta_m$  create an orthogonal system of functions, i.e. the functions satisfy the orthogonality condition

$$\int_{\mu_\varphi}^1 P_{\beta_m}(\mu) P_{\beta_{m'}}(\mu) d\mu = \begin{cases} 0 & \text{for } m' \neq m \\ N_m^\mu & \text{for } m' = m \end{cases} \quad (32)$$

### **Solution of the equation (27)**

The general solution of equation (27) has the form

$$R_{im}(r) = B_{i1} j_{\beta_m}(\omega_{im} r) + B_{i2} y_{\beta_m}(\omega_{im} r), \quad i = 1, 2, \dots, n \quad (33)$$

where  $j_{\beta_m}$  and  $y_{\beta_m}$  are spherical Bessel functions of the first and second kind, respectively. Substituting function  $R_{im}(r)$  into equation (25) and next, the obtained function  $\theta_{im}(r, \mu, t)$  into conditions (13)-(14) and (16)-(17), the system of  $2n$  homogeneous equations is received. We rewrite the equation system in the matrix form

$$\mathbf{A} \cdot \mathbf{B} = \mathbf{0} \quad (34)$$

where  $\mathbf{A} = [a_{ij}]_{1 \leq i, j \leq 2n}$  and  $\mathbf{B} = [B_{11} \ B_{12} \ B_{21} \ B_{22} \ \dots \ B_{n1} \ B_{n2}]^T$ .

The non-zero solutions of equation (34) exist for these values of  $\gamma$  for which the determinant of the matrix  $\mathbf{A}$  disappears, i.e.

$$\det \mathbf{A} = 0 \quad (35)$$

Solving this equation for  $\beta_m$ , the sequences of  $(\gamma_{km})_{k=1,2,\dots}$  are obtained. For each value of  $\omega_{ikm} = \frac{\gamma_{km}}{\sqrt{a_i}}$ , the coefficients  $B_{i1}, B_{i2}$  ( $i = 1, 2, \dots, n$ ) occurring in equation (33) are determined by solving equation (34). The functions  $R_{ikm}(r)$  corresponding to the values of  $\gamma_{km}$  satisfy the orthogonality condition

$$\sum_{i=1}^n \frac{\lambda_i}{a_i} \int_{r_{i-1}}^{r_i} r^2 R_{ikm}(r) R_{ik'm}(r) dr = \begin{cases} 0 & \text{for } k' \neq k \\ N_{km}^r & \text{for } k' = k \end{cases} \quad (36)$$

### **Solution of the equation (28)**

Based on the orthogonality conditions (32), (36) and using (25) in the condition (18), we find the initial condition for the function  $A_{km}$

$$A_{km}(0) = \frac{1}{N_{km}^r N_m^\mu} \sum_{i=1}^n \frac{\lambda_i}{a_i} \int_{r_{i-1}}^{r_i} r^2 R_{ikm}(r) \int_{\mu_\varphi}^1 \Psi_m(\mu) F_i^*(r, \mu) d\mu dr \quad (37)$$

The solution of the fractional differential equation (28) satisfying the initial condition (37) has the form

$$A_{km}(t) = \frac{1}{N_{km}^r N_m^\mu} E_\alpha(-\gamma_k^2 t^\alpha) \sum_{i=1}^n \frac{\lambda_i}{a_i} \int_{r_{i-1}}^{r_i} r^2 R_{ikm}(r) \int_{\mu_\varphi}^1 \Psi_m(\mu) F_i^*(r, \mu) d\mu dr \quad (38)$$

where  $E_{\alpha, \alpha'}(z) = \sum_{k=0}^{\infty} \frac{z^k}{\Gamma(\alpha k + \alpha')}$  is Mittag-Leffler function [12].

Ultimately, using (25), (30), (33) and (38), we have

$$\theta_i(r, \mu, t) = \sum_{k=1}^{\infty} \sum_{m=1}^{\infty} (B_{i1} j_{\beta_m}(\omega_{ikm} r) + B_{i2} y_{\beta_m}(\omega_{ikm} r)) P_{\beta_m}(\mu) A_{km}(t) \quad (39)$$

where  $A_{km}(t)$  is given by equation (38).

### 3.2. Solution to the steady-state problem

We seek a solution of the problem (19)-(24) in the form of a series

$$\Phi_i(r, \mu) = \sum_{m=1}^{\infty} \varphi_{im}(r) P_{\beta_m}(\mu) \quad (40)$$

where  $\beta_m$  are roots of the equation (31a) or (31b). Substituting the function  $\Phi_i(r, \mu)$  into equation (19) we obtain an Euler differential equation

$$r^2 \frac{d^2 \varphi_{im}(r)}{dr^2} + 2r \frac{d\varphi_{im}(r)}{dr} - \beta_m(\beta_m + 1) \varphi_{im}(r) = 0 \quad (41)$$

Next, taking into account function (40) in equations (20)-(21) and (23)-(24) and using the orthogonality condition (32), the boundary conditions for the functions  $\varphi_{im}(r)$  are obtained

$$\lambda_1 \frac{d\varphi_{1m}}{dr}(r_0) - a_+ \varphi_{1m}(r_0) = -\frac{a_+}{N_m^\mu} \int_{\mu_\varphi}^1 T_+(\mu) P_{\beta_m}(\mu) d\mu \quad (42)$$

$$\lambda_n \frac{d\varphi_{nm}}{dr}(r_n) + a_\infty \varphi_{nm}(r_n) = \frac{a_\infty}{N_m^\mu} \int_{\mu_\varphi}^1 T_\infty(\mu) P_{\beta_m}(\mu) d\mu \quad (43)$$

$$\varphi_{im}(r_i) = \varphi_{i+1m}(r_i), \quad i = 1, \dots, n-1 \quad (44)$$

$$\lambda_i \frac{d\varphi_{im}}{dr}(r_i) = \lambda_{i+1} \frac{d\varphi_{i+1m}}{dr}(r_i), \quad i = 1, \dots, n-1 \quad (45)$$

The general solution of the Euler equation (41) has the form

$$\varphi_{im}(r) = C_{i1}r^{\beta_m} + C_{i2}r^{-\beta_m-1} \quad (46)$$

Substitution (46) in conditions (42)-(45) gives  $2n$  linear non-homogeneous equations which can be written in the matrix form

$$\mathbf{D} \cdot \mathbf{C} = \mathbf{F} \quad (47)$$

where  $\mathbf{D} = [d_{ij}]_{1 \leq i, j \leq 2n}$ ,  $\mathbf{C} = [C_{11} \ C_{12} \ \dots \ C_{n1} \ C_{n2}]^T$  and  $\mathbf{F} = \text{col}[f_i]_{1 \leq i \leq 2n}$ .

The equation (47) is then solved with respect to unknown  $\mathbf{C}$ . The determined coefficients  $C_{11}, C_{12}, \dots, C_{n1}, C_{n2}$  are then used in equation (46). Thus, the function  $\Phi_i(r, \mu)$  as a solution of the steady-state problem is given by equation (40) where the functions  $\varphi_{im}(r)$  are defined by equation (46).

Finally, the temperature distribution in the spherical layers received on the basis of the fractional heat conduction model is given by equation (11), (39) and (40).

#### 4. Numerical example

We use the solution of the heat conduction problem derived in Section 3 to numerical calculations of the temperature distribution in a layered hemisphere ( $\mu_\phi = 0$ ). We assume that zero temperature is established at the surface  $\mu = 0$ , i.e. the boundary conditions (15a) and (22a), are satisfied. The considered hemisphere consists of five concentric layers whose locations are determined by non-dimensional radii:  $\bar{r}_i = r_i/r_0$ , where  $r_0 = 1.0$  m. The non-dimensional radii  $\bar{r}_i$ , thermal conductivity  $\lambda_i$  and thermal diffusivity  $a_i$  in the  $i$ -th layer of the sphere are:  $\bar{r}_i = 1.02; 1.04; 1.06; 1.1; 1.2$ ,  $\lambda_i = 16.0; 24.0; 36.0; 54.0; 81.0$  W/(m·K),  $a_i = 3.3 \cdot 10^{-6}; 6.0 \cdot 10^{-6}; 1.1 \cdot 10^{-5}; 2.0 \cdot 10^{-5}; 4.0 \cdot 10^{-5}$  m<sup>2</sup>/s <sup>$\alpha$</sup> ,  $i = 1, \dots, 5$ . The inner and outer heat transfer coefficients are:  $a_+ = a_\infty = 1200$  W/(m<sup>2</sup>·K), the inner and outer ambient temperatures are:  $T_+ = 100.0^\circ\text{C}$ ,  $T_\infty = 20^\circ\text{C}$  and the initial temperature is assumed as constant:  $F_i(r, \mu) = 20^\circ\text{C}$ . The computations were performed for various values of the order of fractional time-derivative:  $\alpha = 0.5; 0.6; 0.7; 0.8; 0.9; 1.0$ . The temperature distributions on the surfaces of the layers determined for  $\mu \in [0, 1]$ , are shown in Figures 2a-d. The temperature depends on the order  $\alpha$  of the time-derivative occurring in the heat conduction model. This dependence is more significant for higher temperature of the sphere.



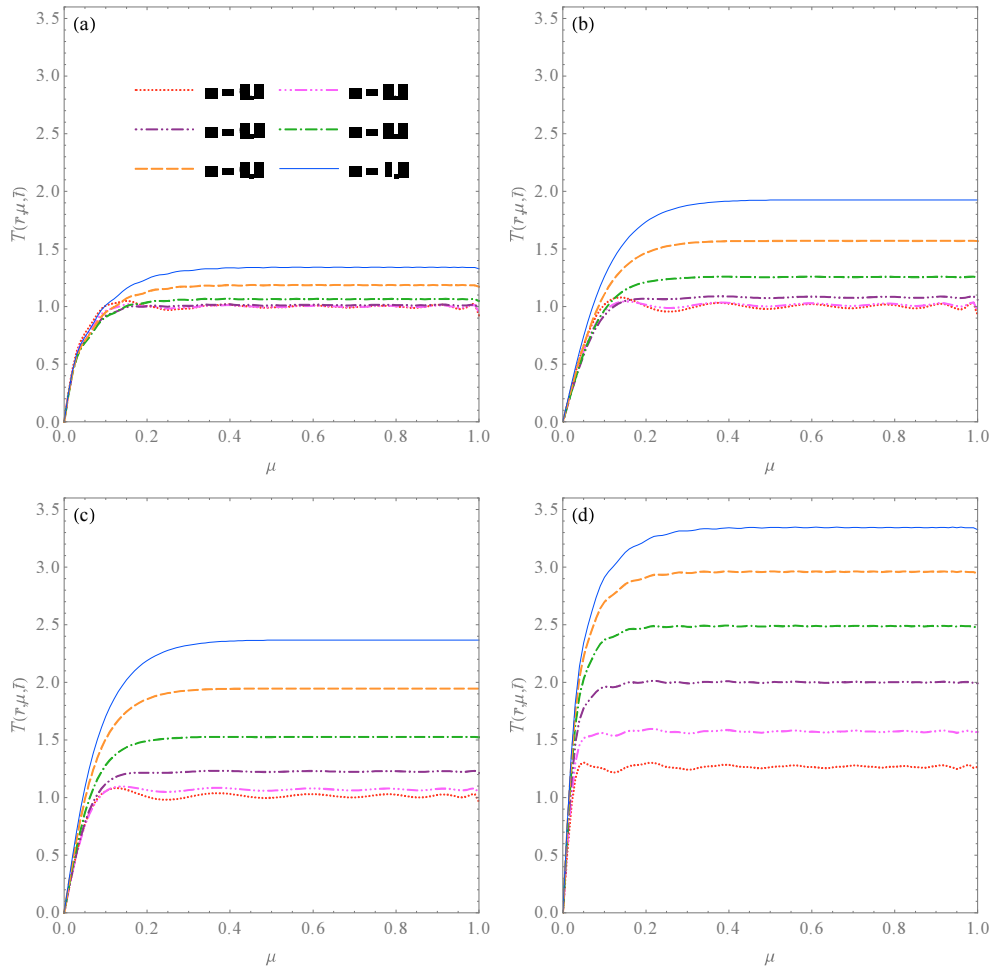


Fig. 2. The non-dimensional temperatures at outer surface and at interfaces of the layered hemisphere: a)  $\bar{r} = 1.2$ , b)  $\bar{r} = 1.1$ , c)  $\bar{r} = 1.06$ , d)  $\bar{r} = 1.02$

## 5. Conclusions

An analytical solution of the problem of time-fractional heat conduction in two-dimensional multi-layer spherical bodies has been presented. The formulation of the problem includes the heat conduction in the spherical bodies which occupied regions defined by finite intervals of the radial coordinate and polar angle. The conditions to perfect contact at interfaces and the mathematical Robin boundary conditions were assumed. The derived solution applies to the bodies consisting of an arbitrary number of layers which are characterized by different thermal conductivity and thermal diffusivity. The numerical example shows that the order of fractional time-derivative is of significant importance for temperature distribution in the body. The temperature at the outer surface and at interfaces of

the layered hemisphere is lower order of the fractional time-derivative. The further research will take into consideration the fractional heat conduction in spherical multilayer bodies with time-dependent boundary conditions.

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