

## 1D GENERALIZED DUAL-PHASE LAG EQUATION. SENSITIVITY ANALYSIS WITH RESPECT TO THE POROSITY

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**Abstract.** Thermal processes occurring in the heated tissue are described by the 1D generalized dual-phase lag equation supplemented by appropriate boundary and initial conditions. Using the sensitivity analysis method, the additional problem connected with the porosity is formulated. Both problems are solved by means of the explicit scheme of the finite difference method. In this way it is possible to estimate the temperature changes due to the perturbation of porosity. In the final part of the paper, the example of computation is shown and the conclusions are formulated.

**Keywords:** *bioheat transfer, generalized dual-phase lag model, sensitivity analysis, finite difference method*

### 1. Introduction

Thermal phenomena occurring in heated tissue are described using the various mathematical models, e.g. [1-8]. One of them is the dual-phase lag equation [9-11] in which two time parameters  $\tau_q$  and  $\tau_T$  appear. The parameter  $\tau_q$  is the phase lag of the heat flux and parameter  $\tau_T$  is the phase lag of the temperature gradient. Extension of this model is the generalized dual-phase lag equation [12-15]. In this equation the phase lag times  $\tau_q$  and  $\tau_T$  are expressed in terms of the blood and tissue properties, the interphase convective heat transfer coefficient and the blood perfusion rate. Generalized dual-phase lag equation contains additional parameters, it means the coupling factor  $G$ , the porosity  $\varepsilon$  defined as a ratio of blood volume to total volume and the blood temperature  $T_b$ .

In this paper the 1D generalized dual-phase lag equation (GDPLE) under the assumption that the blood temperature is constant is considered. The aim of investigations is to estimate the changes of the temperature field due to the perturbations of porosity  $\varepsilon$ . For this purpose, the direct method of sensitivity analysis is applied [16-20]. In this case the mathematical model consists of the basic problem (GDPLE supplemented by appropriate boundary and initial conditions) and the

additional problem resulting from the differentiation of governing equations. These problems are solved using the explicit scheme of the finite difference method for hyperbolic equations [14, 15]. In the final part of the paper the example of computations and conclusions are presented.

## 2. Formulation of the problem

The generalized dual-phase lag equation describing the heat transfer processes occurring in the heated tissues is based on the theory of porous media [6, 12, 13]. In this approach the tissue is divided into two regions: the vascular region (blood vessels) and the extravascular region (tissue). An important parameter in this model is the porosity  $\varepsilon$ , it means the ratio of blood volume to total volume. Using this parameter, the effective thermal conductivity

$$\Lambda = (1 - \varepsilon)\lambda + \varepsilon\lambda_b \quad (1)$$

and effective heat capacity

$$C = (1 - \varepsilon)\rho c + \varepsilon\rho_b c_b \quad (2)$$

are defined, where  $\lambda$ ,  $\lambda_b$  are the thermal conductivities of tissue and blood, respectively,  $\rho$ ,  $\rho_b$  are the densities,  $c$ ,  $c_b$  are the specific heats.

Under the assumption that the metabolic heat sources  $Q$ ,  $Q_b$  and the blood temperature  $T_b$  are constant, the generalized dual-phase lag equation for 1D problem takes the form [13, 15, 16]

$$0 < x < L: \quad C \frac{\partial T}{\partial t} + C\tau_q \frac{\partial^2 T}{\partial t^2} = \Lambda \frac{\partial^2 T}{\partial x^2} + \Lambda\tau_T \frac{\partial^3 T}{\partial t \partial x^2} + G(T_b - T) + \varepsilon Q_b + (1 - \varepsilon)Q \quad (3)$$

where  $T = T(x, t)$  is the temperature,  $x$  is the spatial co-ordinate,  $t$  is the time,  $G$  is the coupling factor,  $\tau_q$  and  $\tau_T$  are the phase lags for heat flux and temperature gradient. The equation (3) is supplemented by boundary conditions [10, 11]

$$x = 0: \quad -\Lambda \frac{\partial T}{\partial x} - \Lambda\tau_T \frac{\partial^2 T}{\partial t \partial x} = q_b \quad (4)$$

$$x = L: \quad T = T_b$$

and the initial ones

$$t = 0: \quad T = T_p, \quad \frac{\partial T}{\partial t} = 0 \quad (5)$$

where  $q_b$  is the known constant boundary heat flux,  $T_b$  is the boundary temperature and  $T_p$  is the initial temperature of biological tissue.

The coupling factor  $G$  and the phase lags  $\tau_q, \tau_T$  are defined as [12, 14]

$$G = A\alpha + w c_b \quad (6)$$

$$\tau_q = \frac{\varepsilon(1-\varepsilon)\rho c \rho_b c_b}{GC} \quad (7)$$

$$\tau_T = \frac{\varepsilon(1-\varepsilon)\rho_b c_b \lambda}{G\Lambda} \quad (8)$$

where  $A$  is the the volumetric heat transfer area between tissue and blood,  $\alpha$  is the heat transfer coefficient and  $w$  is the blood perfusion rate.

The volumetric heat transfer area between tissue and blood can be determined from the dependence  $A = k P_b/V$ , where  $k$  is the number of blood vessels located in the volume  $V$  and  $P_b$  is the surface area between the single vessel and tissue. As mentioned before, the porosity  $\varepsilon$  is the ratio of blood volume to total volume:  $\varepsilon = kV_b/V$  and it follows that  $A = \varepsilon P_b/V_b$ . The heat transfer coefficient can be calculated from the Nusselt number  $Nu = \alpha d/\lambda_b$ , where  $d$  is the vessel diameter. Assuming that the diameter of all blood vessels is the same, then  $P_b/V_b$  and  $\alpha$  are constants.

### 3. Sensitivity analysis

To estimate the temperature changes due to the changes of porosity, the direct method of sensitivity analysis can be applied [16-20].

The equation (3) is differentiated with respect to the porosity  $\varepsilon$  and then

$$\begin{aligned} \frac{\partial C}{\partial \varepsilon} \frac{\partial T}{\partial t} + C \frac{\partial}{\partial t} \left( \frac{\partial T}{\partial \varepsilon} \right) + \frac{\partial (C\tau_q)}{\partial \varepsilon} \frac{\partial^2 T}{\partial t^2} + C\tau_q \frac{\partial^2}{\partial t^2} \left( \frac{\partial T}{\partial \varepsilon} \right) = \frac{\partial \Lambda}{\partial \varepsilon} \frac{\partial^2 T}{\partial x^2} + \Lambda \frac{\partial^2}{\partial x^2} \left( \frac{\partial T}{\partial \varepsilon} \right) + \\ \frac{\partial (\Lambda\tau_T)}{\partial \varepsilon} \frac{\partial^3 T}{\partial t \partial x^2} + \Lambda\tau_T \frac{\partial^3}{\partial t \partial x^2} \left( \frac{\partial T}{\partial \varepsilon} \right) + \frac{\partial G}{\partial \varepsilon} (T_b - T) - G \frac{\partial T}{\partial \varepsilon} + Q_b - Q \end{aligned} \quad (9)$$

or

$$\begin{aligned} C \frac{\partial U}{\partial t} + C\tau_q \frac{\partial^2 U}{\partial t^2} = \Lambda \frac{\partial^2 U}{\partial x^2} + \Lambda\tau_T \frac{\partial^3 U}{\partial t \partial x^2} + \frac{\partial \Lambda}{\partial \varepsilon} \frac{\partial^2 T}{\partial x^2} + \frac{\partial (\Lambda\tau_T)}{\partial \varepsilon} \frac{\partial^3 T}{\partial t \partial x^2} - \\ \frac{\partial C}{\partial \varepsilon} \frac{\partial T}{\partial t} - \frac{\partial (C\tau_q)}{\partial \varepsilon} \frac{\partial^2 T}{\partial t^2} + \frac{\partial G}{\partial \varepsilon} (T_b - T) - GU + Q_b - Q \end{aligned} \quad (10)$$

where  $U = \partial T / \partial \varepsilon$  is the sensitivity function.

The boundary conditions (4) and the initial conditions (5) are also differentiated with respect to  $\varepsilon$ . Thus

$$\begin{aligned} x=0: & -\frac{\partial \Lambda}{\partial \varepsilon} \frac{\partial T}{\partial x} - \Lambda \frac{\partial U}{\partial x} - \frac{\partial(\Lambda \tau_T)}{\partial \varepsilon} \frac{\partial^2 T}{\partial t \partial x} - \Lambda \tau_T \frac{\partial^2 U}{\partial t \partial x} = 0 \\ x=L: & U=0 \end{aligned} \quad (11)$$

and

$$t=0: U=0, \quad \frac{\partial U}{\partial t} = 0 \quad (12)$$

In equation (10) the derivatives of thermal parameters are determined as follows

$$\frac{\partial G}{\partial \varepsilon} = \frac{P_b}{V_b} \alpha, \quad \frac{\partial \Lambda}{\partial \varepsilon} = \lambda_b - \lambda, \quad \frac{\partial C}{\partial \varepsilon} = \rho_b c_b - \rho c \quad (13)$$

$$\frac{\partial(\Lambda \tau_T)}{\partial \varepsilon} = \frac{1}{G} (1-2\varepsilon) \rho_b c_b \lambda - \frac{1}{G^2} \frac{\partial G}{\partial \varepsilon} \varepsilon (1-\varepsilon) \rho_b c_b \lambda \quad (14)$$

$$\frac{\partial(C \tau_q)}{\partial \varepsilon} = \frac{1}{G} (1-2\varepsilon) \rho c \rho_b c_b - \frac{1}{G^2} \frac{\partial G}{\partial \varepsilon} \varepsilon (1-\varepsilon) \rho c \rho_b c_b \quad (15)$$

The equations (10)-(12) create the additional problem which is coupled with the basic problem (equations (3)-(5)).

#### 4. Method of solution

To solve both the basic problem and the additional one, the finite difference method in the version presented in [21] is applied. The geometrical mesh with constant step  $h$  and time step  $\Delta t$  are introduced.

For the internal node  $i$  the following approximate form of equation (3) is proposed

$$\begin{aligned} C \frac{T_i^f - T_i^{f-1}}{\Delta t} + C \tau_q \frac{T_i^f - 2T_i^{f-1} + T_i^{f-2}}{(\Delta t)^2} &= \frac{\Lambda}{h^2} (T_{i-1}^{f-1} - 2T_i^{f-1} + T_{i+1}^{f-1}) + \\ \frac{\Lambda \tau_T}{h^2 \Delta t} \left[ (T_{i-1}^{f-1} - 2T_i^{f-1} + T_{i+1}^{f-1}) - (T_{i-1}^{f-2} - 2T_i^{f-2} + T_{i+1}^{f-2}) \right] &+ \\ G(T_b - T_i^{f-1}) + \varepsilon Q_b + (1-\varepsilon) Q & \end{aligned} \quad (16)$$

After mathematical manipulations one obtains ( $i = 1, 2, \dots, n - 1$ )

$$T_i^f = \left[ A_3 C (\Delta t + 2\tau_q) - 2A_1 - A_3 (\Delta t)^2 G \right] T_i^{f-1} - (A_3 C \tau_q - 2A_2) T_i^{f-2} + A_1 (T_{i-1}^{f-1} + T_{i+1}^{f-1}) - A_2 (T_{i-1}^{f-2} + T_{i+1}^{f-2}) + A_3 (\Delta t)^2 G T_b + A_3 (\Delta t)^2 [\varepsilon Q_b + (1 - \varepsilon) Q] \quad (17)$$

where

$$A_1 = \frac{\Lambda \Delta t (\Delta t + \tau_T)}{Ch^2 (\Delta t + \tau_q)}, \quad A_2 = \frac{\Lambda \tau_T \Delta t}{Ch^2 (\Delta t + \tau_q)}, \quad A_3 = \frac{1}{C (\Delta t + \tau_q)} \quad (18)$$

For the boundary node  $i = 0$  one has

$$-\Lambda \left( \frac{\partial T}{\partial x} \right)_0^f - \frac{\Lambda \tau_T}{\Delta t} \left[ \left( \frac{\partial T}{\partial x} \right)_0^f - \left( \frac{\partial T}{\partial x} \right)_0^{f-1} \right] = q_b \quad (19)$$

and next

$$-\frac{\Lambda (\Delta t + \tau_T)}{h \Delta t} (T_1^f - T_0^f) + \frac{\Lambda \tau_T}{h \Delta t} (T_1^{f-1} - T_0^{f-1}) = q_b \quad (20)$$

From equation (20) it follows that

$$T_0^f = T_1^f - \frac{\tau_T}{\Delta t + \tau_T} (T_1^{f-1} - T_0^{f-1}) + \frac{h \Delta t}{\Lambda (\Delta t + \tau_T)} q_b \quad (21)$$

In the similar way the equation (10) connected with the sensitivity function is approximated

$$\begin{aligned} & \frac{C (\Delta t + \tau_q)}{(\Delta t)^2} U_i^f - \frac{C (\Delta t + 2\tau_q)}{(\Delta t)^2} U_i^{f-1} + \frac{C \tau_q}{(\Delta t)^2} U_i^{f-2} = \\ & \frac{\Lambda (\Delta t + \tau_T)}{h^2 \Delta t} (U_{i-1}^{f-1} - 2U_i^{f-1} + U_{i+1}^{f-1}) - \frac{\Lambda \tau_T}{h^2 \Delta t} (U_{i-1}^{f-2} - 2U_i^{f-2} + U_{i+1}^{f-2}) - G U_i^{f-1} + \\ & \frac{1}{h^2 \Delta t} \left( \Delta t \frac{\partial \Lambda}{\partial \varepsilon} + \frac{\partial (\Lambda \tau_T)}{\partial \varepsilon} \right) (T_{i-1}^f - 2T_i^f + T_{i+1}^f) - \frac{1}{h^2 \Delta t} \frac{\partial (\Lambda \tau_T)}{\partial \varepsilon} (T_{i-1}^{f-1} - 2T_i^{f-1} + T_{i+1}^{f-1}) - \\ & \frac{1}{(\Delta t)^2} \left( \Delta t \frac{\partial C}{\partial \varepsilon} + \frac{\partial (C \tau_q)}{\partial \varepsilon} + (\Delta t)^2 \frac{\partial G}{\partial \varepsilon} \right) T_i^f + \frac{1}{(\Delta t)^2} \left( \Delta t \frac{\partial C}{\partial \varepsilon} + 2 \frac{\partial (C \tau_q)}{\partial \varepsilon} \right) T_i^{f-1} - \\ & \frac{1}{(\Delta t)^2} \frac{\partial (C \tau_q)}{\partial \varepsilon} T_i^{f-2} + \frac{\partial G}{\partial \varepsilon} T_b + Q_b - Q \end{aligned} \quad (22)$$

From the equation (22) it results that ( $i = 1, 2, \dots, n - 1$ )

$$\begin{aligned}
U_i^f = & \left[ A_3 C (\Delta t + 2\tau_q) - 2A_1 - A_3 G (\Delta t)^2 \right] U_i^{f-1} - \\
& \left[ A_3 C \tau_q - 2A_2 \right] U_i^{f-2} + A_1 (U_{i-1}^{f-1} + U_{i+1}^{f-1}) - A_2 (U_{i-1}^{f-2} + U_{i+1}^{f-2}) + \\
& A_3 \frac{\Delta t}{h^2} \left( \Delta t \frac{\partial \Lambda}{\partial \varepsilon} + \frac{\partial (\Lambda \tau_T)}{\partial \varepsilon} \right) (T_{i-1}^f - 2T_i^f + T_{i+1}^f) - A_3 \frac{\Delta t}{h^2} \frac{\partial (\Lambda \tau_T)}{\partial \varepsilon} (T_{i-1}^{f-1} - 2T_i^{f-1} + T_{i+1}^{f-1}) - \quad (23) \\
& A_3 \left( \Delta t \frac{\partial C}{\partial \varepsilon} + \frac{\partial (C \tau_q)}{\partial \varepsilon} + (\Delta t)^2 \frac{\partial G}{\partial \varepsilon} \right) T_i^f + A_3 \left( \Delta t \frac{\partial C}{\partial \varepsilon} + 2 \frac{\partial (C \tau_q)}{\partial \varepsilon} \right) T_i^{f-1} - \\
& A_3 \frac{\partial (C \tau_q)}{\partial \varepsilon} T_i^{f-2} + A_3 (\Delta t)^2 \left( \frac{\partial G}{\partial \varepsilon} T_b + Q_b - Q \right)
\end{aligned}$$

For the boundary node  $i = 0$  one obtains (cf. equation (11))

$$\begin{aligned}
& - \frac{\Lambda (\Delta t + \tau_T)}{h \Delta t} (U_1^f - U_0^f) + \frac{\Lambda \tau_T}{h \Delta t} (U_1^{f-1} - U_0^{f-1}) - \frac{1}{h} \frac{\partial \Lambda}{\partial \varepsilon} (T_1^f - T_0^f) - \\
& \frac{1}{h \Delta t} \frac{\partial (\Lambda \tau_T)}{\partial \varepsilon} [(T_1^f - T_0^f) - (T_1^{f-1} - T_0^{f-1})] = 0 \quad (24)
\end{aligned}$$

hence

$$\begin{aligned}
U_0^f = & U_1^f - \frac{\tau_T}{\Delta t + \tau_T} (U_1^{f-1} - U_0^{f-1}) + \frac{\Delta t}{\Lambda (\Delta t + \tau_T)} \frac{\partial \Lambda}{\partial \varepsilon} (T_1^f - T_0^f) + \\
& \frac{1}{\Lambda (\Delta t + \tau_T)} \frac{\partial (\Lambda \tau_T)}{\partial \varepsilon} [(T_1^f - T_0^f) - (T_1^{f-1} - T_0^{f-1})] \quad (25)
\end{aligned}$$

Because the explicit scheme of the finite difference method is applied, the stability criteria should be fulfilled [21]

$$\begin{aligned}
& A_3 C (\Delta t + 2\tau_q) - 2A_1 - A_3 (\Delta t)^2 G \geq 0 \\
& A_3 C \tau_q - 2A_2 \geq 0 \quad (26)
\end{aligned}$$

## 5. Results of computations

The layer of biological tissue of thickness  $L = 0.01$  m is considered. On the surface  $x = 0$  the Neumann condition  $q_b = 1500$  W/m<sup>2</sup> is assumed and on the surface  $x = L$  the Dirichlet condition  $T_b = 37^\circ\text{C}$  is accepted. The initial temperature of tissue

is equal to  $T_p = 37^\circ\text{C}$ . The thermophysical parameters of tissue and blood are collected in Table 1. The blood perfusion rate  $w = 5 \text{ kg}/(\text{m}^3\text{s})$ , the porosity  $\varepsilon = 0.1637$ , the vessels diameter  $d = 0.00456 \text{ m}$ , the Nusselt number  $\text{Nu} = 4.93$ .

The calculations were made using the explicit scheme of the finite difference method assuming the grid step  $h = 0.0002 \text{ m}$  and time step  $\Delta t = 0.01 \text{ s}$ .

Table 1

Thermophysical parameters

Parameter	Tissue	Blood
Thermal conductivity [W/(m·K)]	0.3	0.5
Density [kg/m <sup>3</sup> ]	1000	1060
Specific heat [J/(kg·K)]	4000	3770
Metabolic heat source [W/m <sup>3</sup> ]	240	250

In Figure 1 the distribution of sensitivity function after 60, 120 and 180 seconds is presented. Figures 2 and 3 illustrate the courses of temperature and sensitivity function at the points  $x = 0$  (heated surface),  $x = 1 \text{ mm}$ ,  $x = 2 \text{ mm}$ ,  $x = 3 \text{ mm}$  and  $x = 4 \text{ mm}$ . As might be expected, the biggest changes of the sensitivity function occur near the heated surface (Fig. 1).

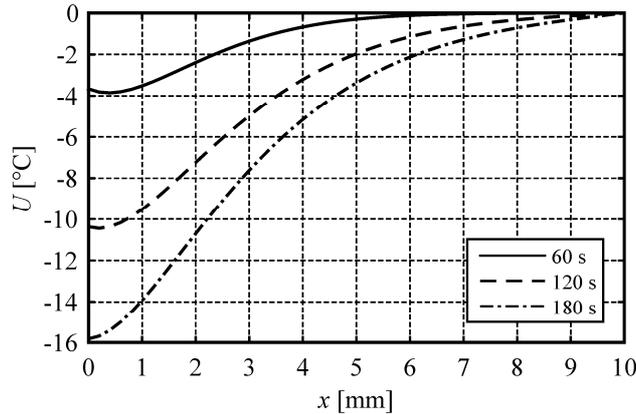


Fig. 1. Distribution of the sensitivity function after 60, 120 and 180 seconds

Knowledge of the sensitivity function allows one, among others, to estimate the changes of temperature due to the perturbation of the parameter considered:  $\Delta T = U\Delta\varepsilon$ , where  $\Delta\varepsilon$  is the assumed perturbation of the porosity (e.g.  $\Delta\varepsilon = 0.05\varepsilon$ ). In Table 2 the changes of temperature for different values of  $\Delta\varepsilon$  at the point  $x = 0$  are collected. As can be seen, changes in temperature increase both with an increase of  $\Delta\varepsilon$  as well as with the duration of the heating process.

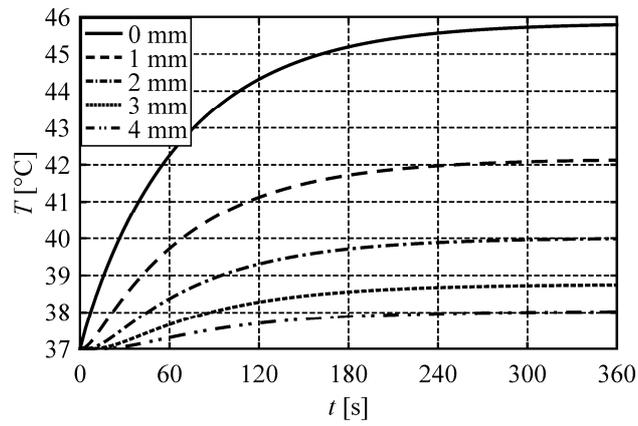


Fig. 2. Heating curves at selected points

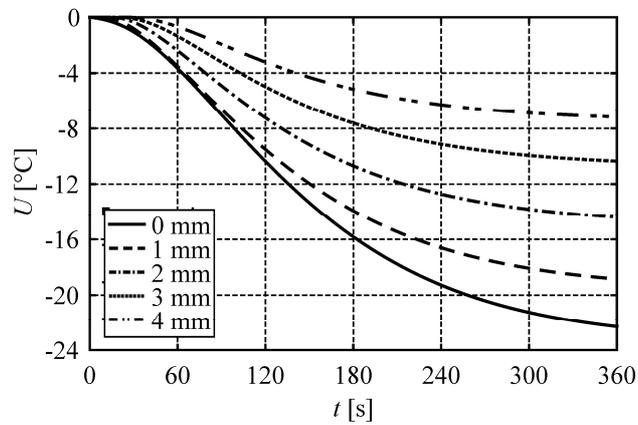


Fig. 3. Course of sensitivity function at selected points

Table 2

Change of temperature due to the perturbation of porosity ( $x = 0$ )

Time [s]	$\Delta T$ [°C]		
	$\Delta\varepsilon = 0.01\varepsilon$	$\Delta\varepsilon = 0.05\varepsilon$	$\Delta\varepsilon = 0.1\varepsilon$
60	-0.01	-0.03	-0.06
120	-0.02	-0.09	-0.17
180	-0.03	-0.13	-0.26
240	-0.03	-0.16	-0.32
300	-0.03	-0.17	-0.35
360	-0.04	-0.18	-0.37

## 6. Conclusions

The 1D generalized dual-phase lag equation describing the thermal processes occurring in the heated biological tissue is considered. The sensitivity analysis with respect to the porosity is presented. Here, the diameter of vessels is assumed as a constant and then the porosity contains information only about the density of blood vessels.

As might be expected, the highest temperature perturbation occurs near the heated boundary, which can be seen in Figure 1. On the other hand, however, the heating time also has a meaningful impact on the results (cf. Fig. 3). It should be noted that in each case, the sensitivity function takes the negative values. This results from an increase in porosity, which entails in decreasing the tissue temperature, of course. As can be seen in Table 2, the porosity perturbation of 1% causes very small changes in temperature ( $< 0.1\%$ ), while the porosity perturbation of 5 or 10% noticeably changes the values of temperature.

The presented approach can be extended for 3D problems and two-temperature models [12, 14, 22] in which the blood temperature is determined from the additional equation coupled with the generalized dual-phase lag equation.

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